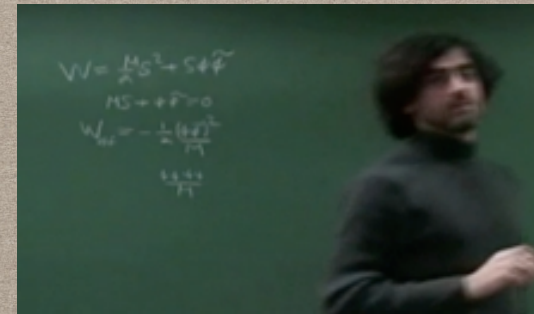


DYNAMICS AT A QUANTUM CRITICAL POINT: QMC + HOLOGRAPHY

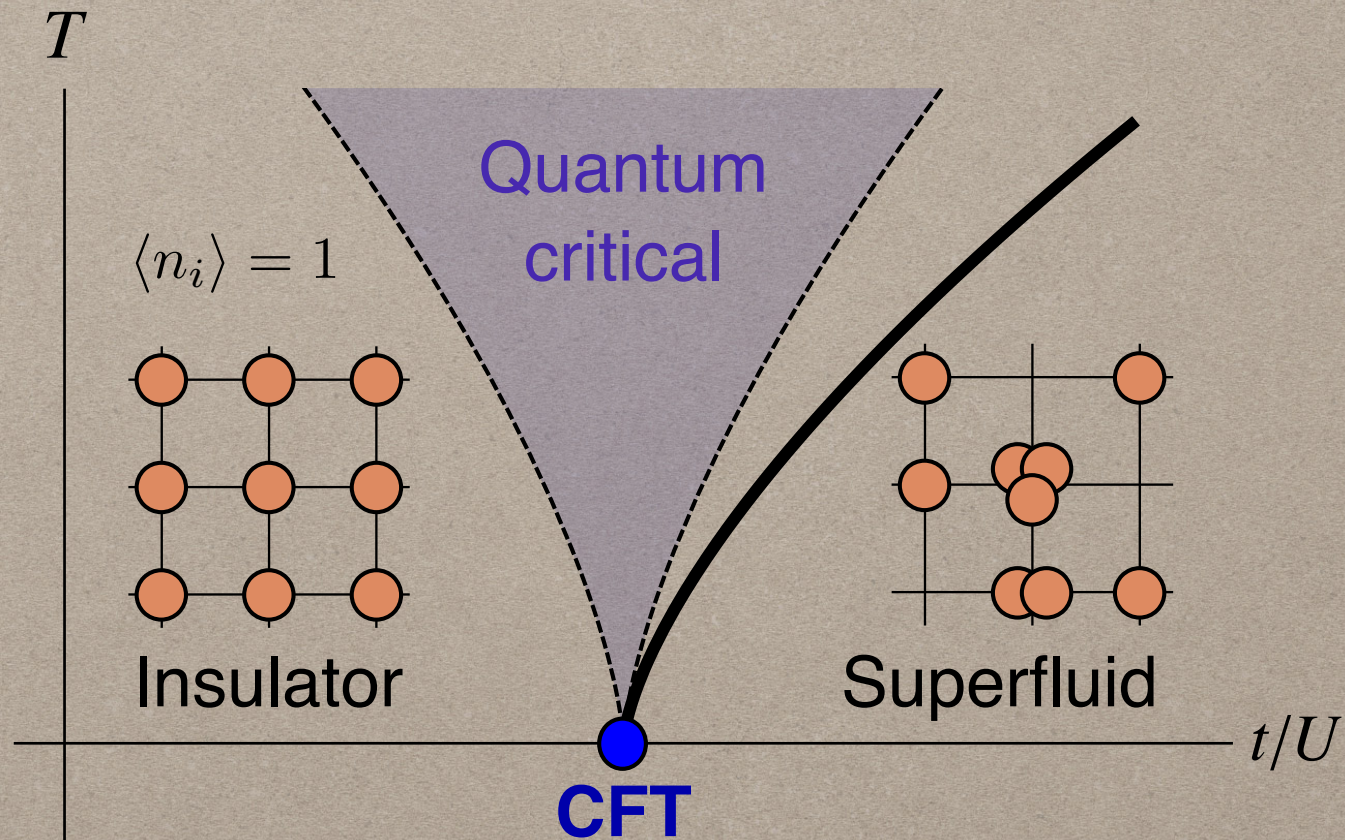
ERIK S SØRENSEN

WILLIAM WITCZAK-KREMPA
SUBIR SACHDEV
EMANUEL KATZ

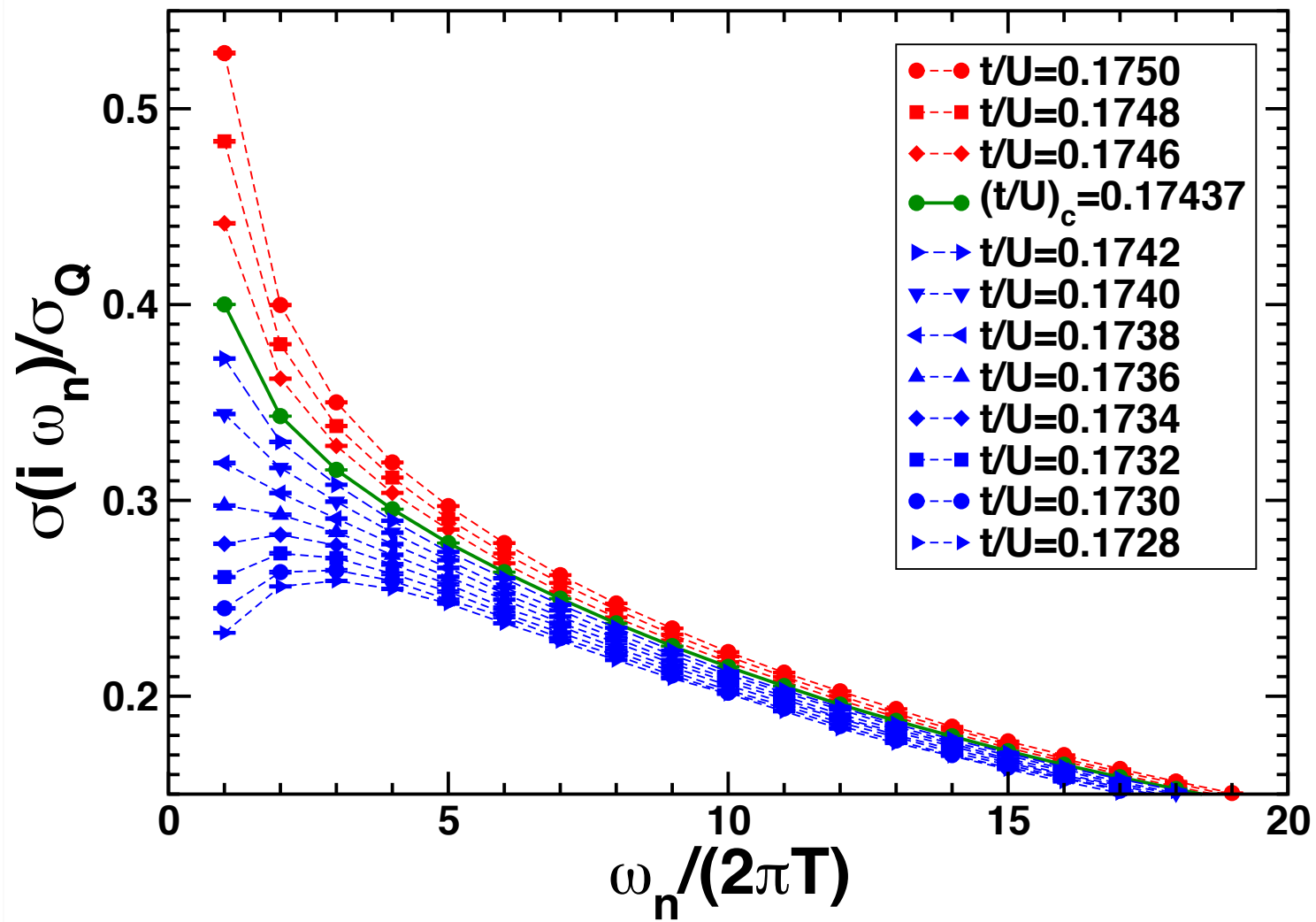


Bose-Hubbard Model in 2+1 D

$$H = -t \sum_{\langle i,j \rangle} b_i^\dagger b_j + U \sum_i n_i(n_i - 1) - \mu \sum_i n_i$$



- $Z=1$ Emergent Lorentz Symmetry
- $2+1D$ $O(2)$ Universality quantum Rotors

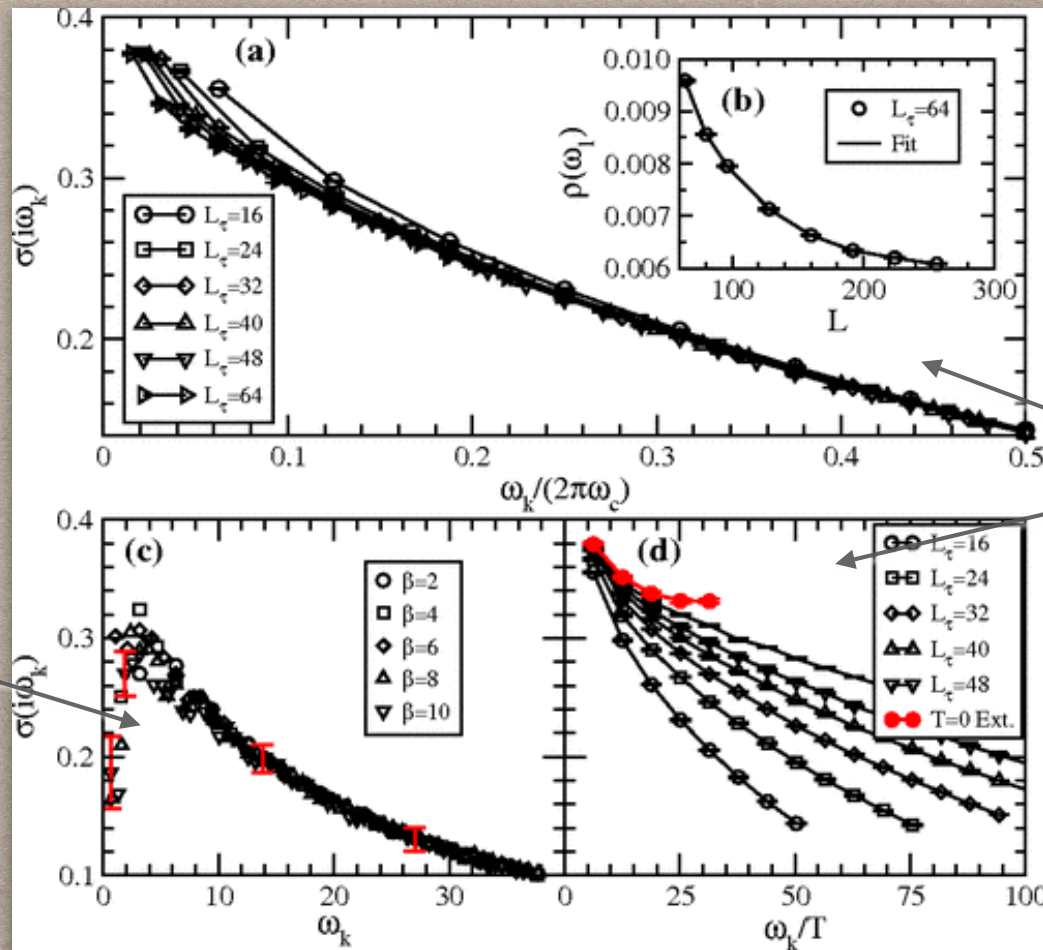


Conductivity: $T^{(d-2)/z} \sigma(\omega, T)$

Scaling:

$$\sigma(\omega/T, T \rightarrow 0) = (k_B T / \hbar c)^{(d-2)/z} \sigma_Q \Sigma(\hbar \omega / k_B T)$$

$$\sigma_Q = (e^*)^2 / h,$$

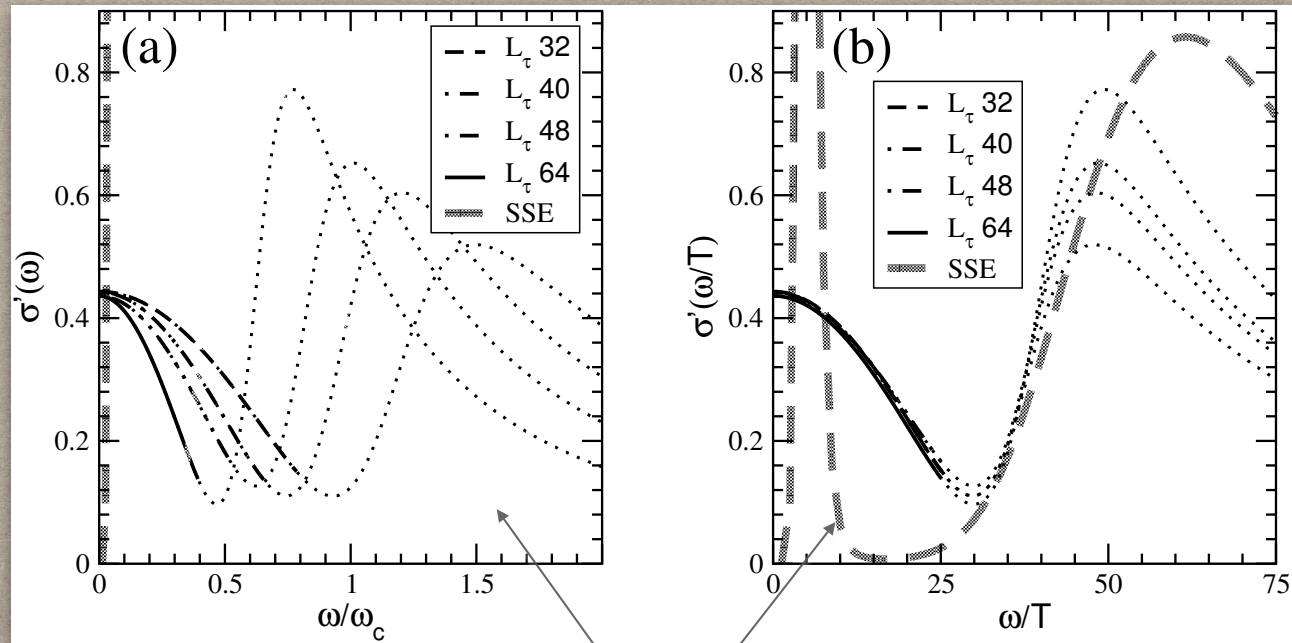


Bose-Hubbard
Model

Villain Model

J. Smakov, ESS PRL 2005

Analytical Continuation



Max Ent

J. Smakov, ESS PRL 2005

Quantum Rotors

$$H_{\text{qr}} = \frac{U}{2} \sum_{\mathbf{r}} \frac{1}{2} \left(\frac{1}{i} \frac{\partial}{\partial \theta_{\mathbf{r}}} \right)^2 - \mu \sum_{\mathbf{r}} \frac{1}{i} \frac{\partial}{\partial \theta_{\mathbf{r}}} - \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} t \cos(\theta_{\mathbf{r}} - \theta_{\mathbf{r}'}).$$

$$Z_{\text{QR}} \approx \sum'_{\{\mathbf{J}\}} \exp \left\{ - \sum_{(\mathbf{r}, \tau)} \left(\Delta\tau U \left(\frac{1}{2} [J_{(\mathbf{r}, \tau)}^{\tau}]^2 - \frac{\mu}{U} J_{(\mathbf{r}, \tau)}^{\tau} \right) - \ln \left(I_{J_{(\mathbf{r}, \tau)}^x}(t\Delta\tau) \right) - \ln \left(I_{J_{(\mathbf{r}, \tau)}^y}(t\Delta\tau) \right) \right) \right\}.$$

$\Delta\tau \rightarrow 0$

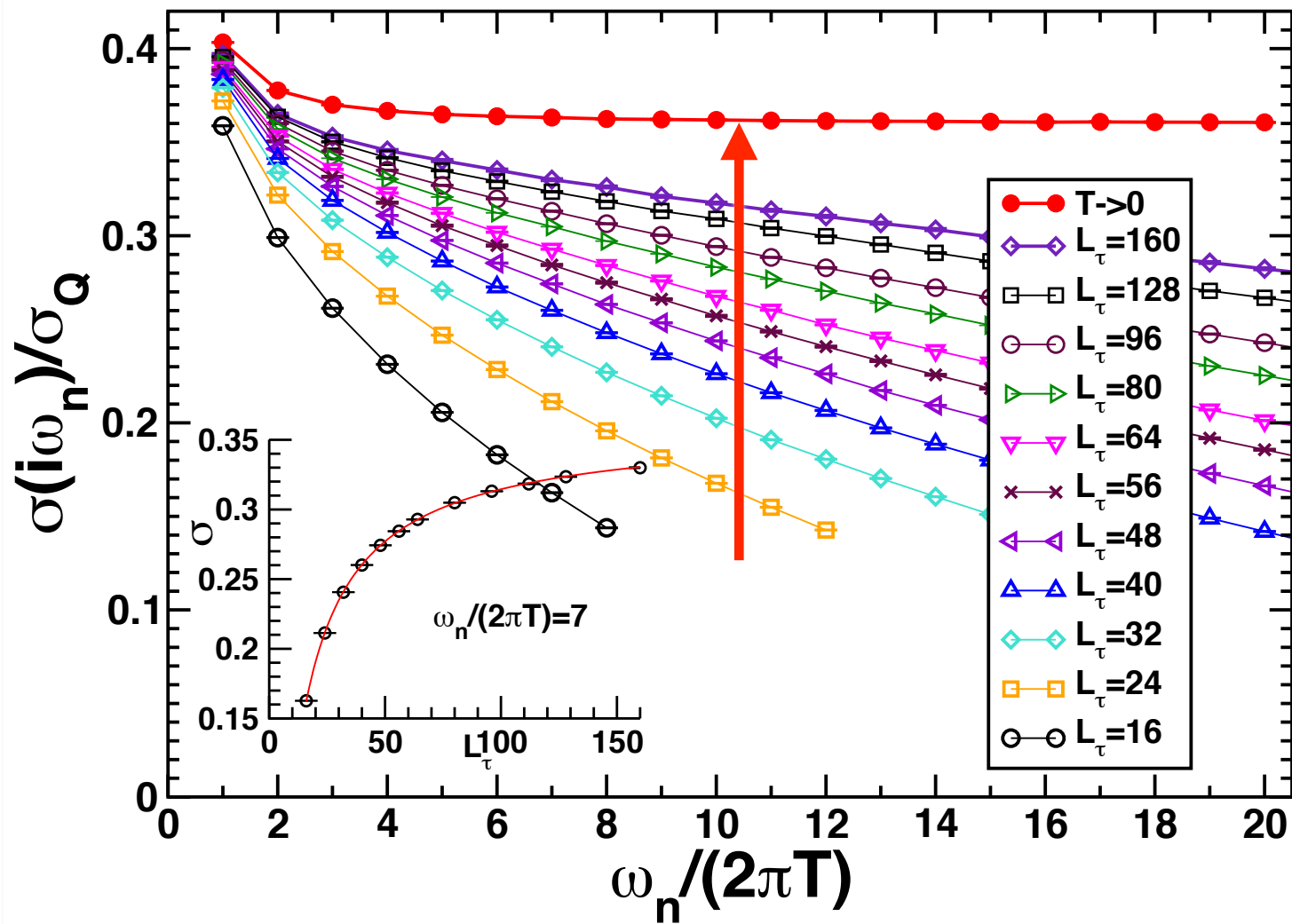
Villain Model

$$H_{\text{V}} = \frac{U}{2} \sum_{\mathbf{r}} \frac{1}{2} \left(\frac{1}{i} \frac{\partial}{\partial \theta_{\mathbf{r}}} \right)^2 - \mu \sum_{\mathbf{r}} \frac{1}{i} \frac{\partial}{\partial \theta_{\mathbf{r}}} - \sum_{\langle \mathbf{r}, \mathbf{r}' \rangle} t \sum_m (\theta_{\mathbf{r}} - \theta_{\mathbf{r}'} - 2\pi m)^2.$$

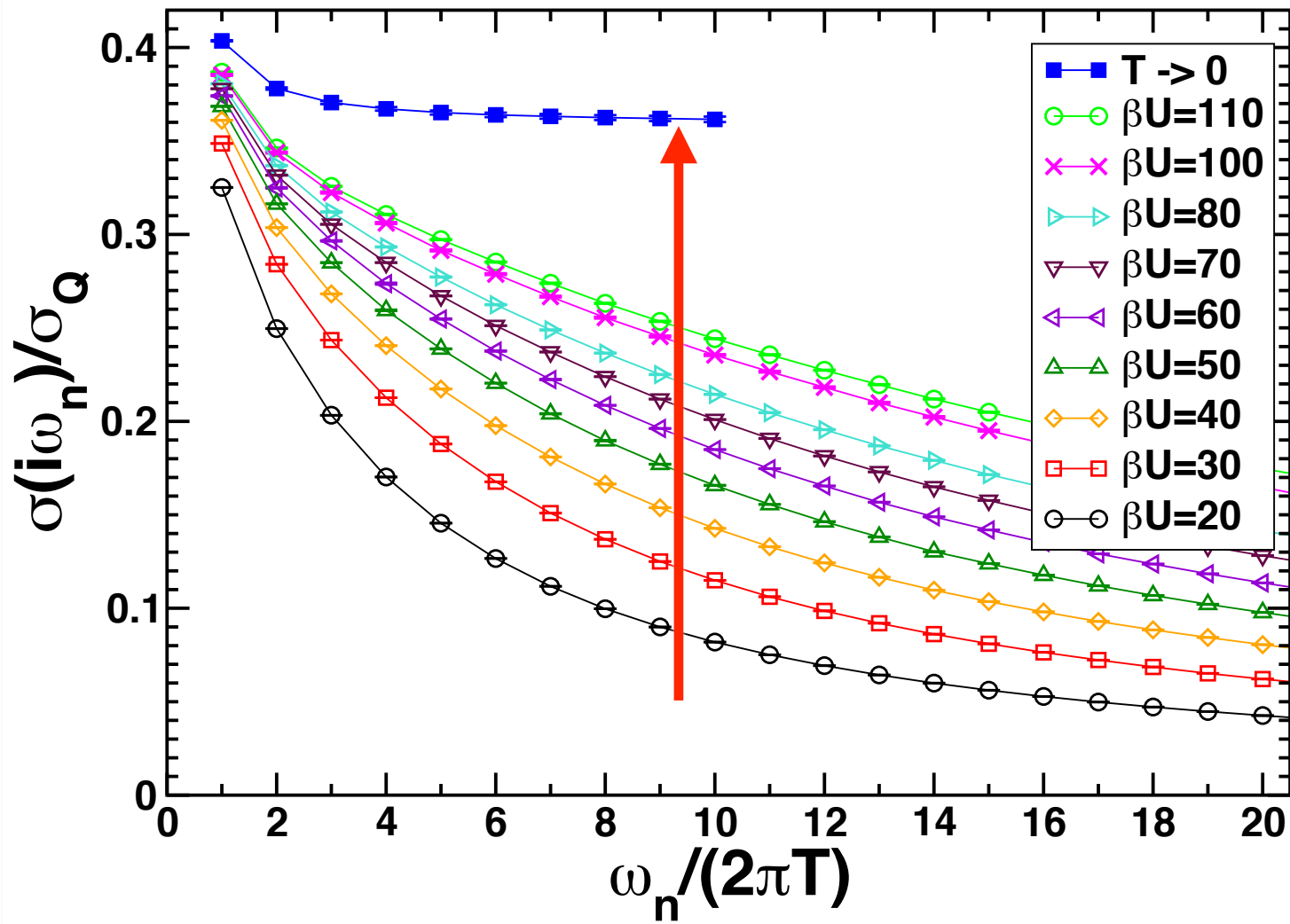
$$Z_{\text{V}} \approx \sum'_{\{\mathbf{J}\}} \exp \left\{ - \frac{1}{K} \sum_{(\mathbf{r}, \tau)} \left(\frac{1}{2} \left([J_{(\mathbf{r}, \tau)}^x]^2 + [J_{(\mathbf{r}, \tau)}^y]^2 + [J_{(\mathbf{r}, \tau)}^{\tau}]^2 \right) - \frac{\mu}{U} J_{(\mathbf{r}, \tau)}^{\tau} \right) \right\}.$$

$$\Delta\tau = \frac{1}{\sqrt{Ut}} \text{ Fixed}$$

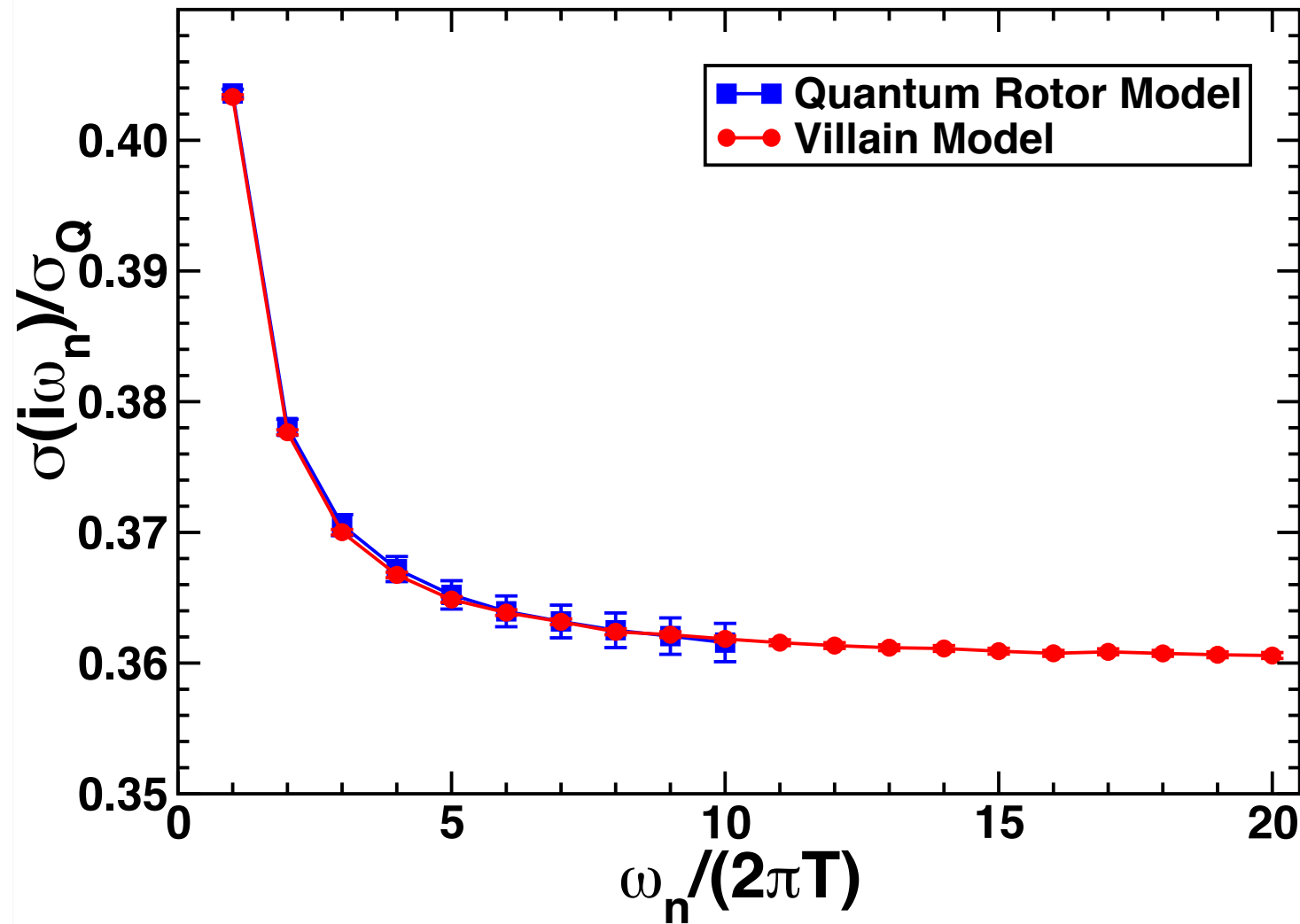
Villain Model



Quantum Rotors

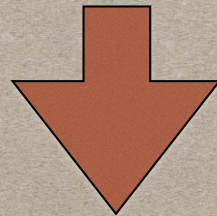


$T \rightarrow 0$ Extrapolated Results



OPE

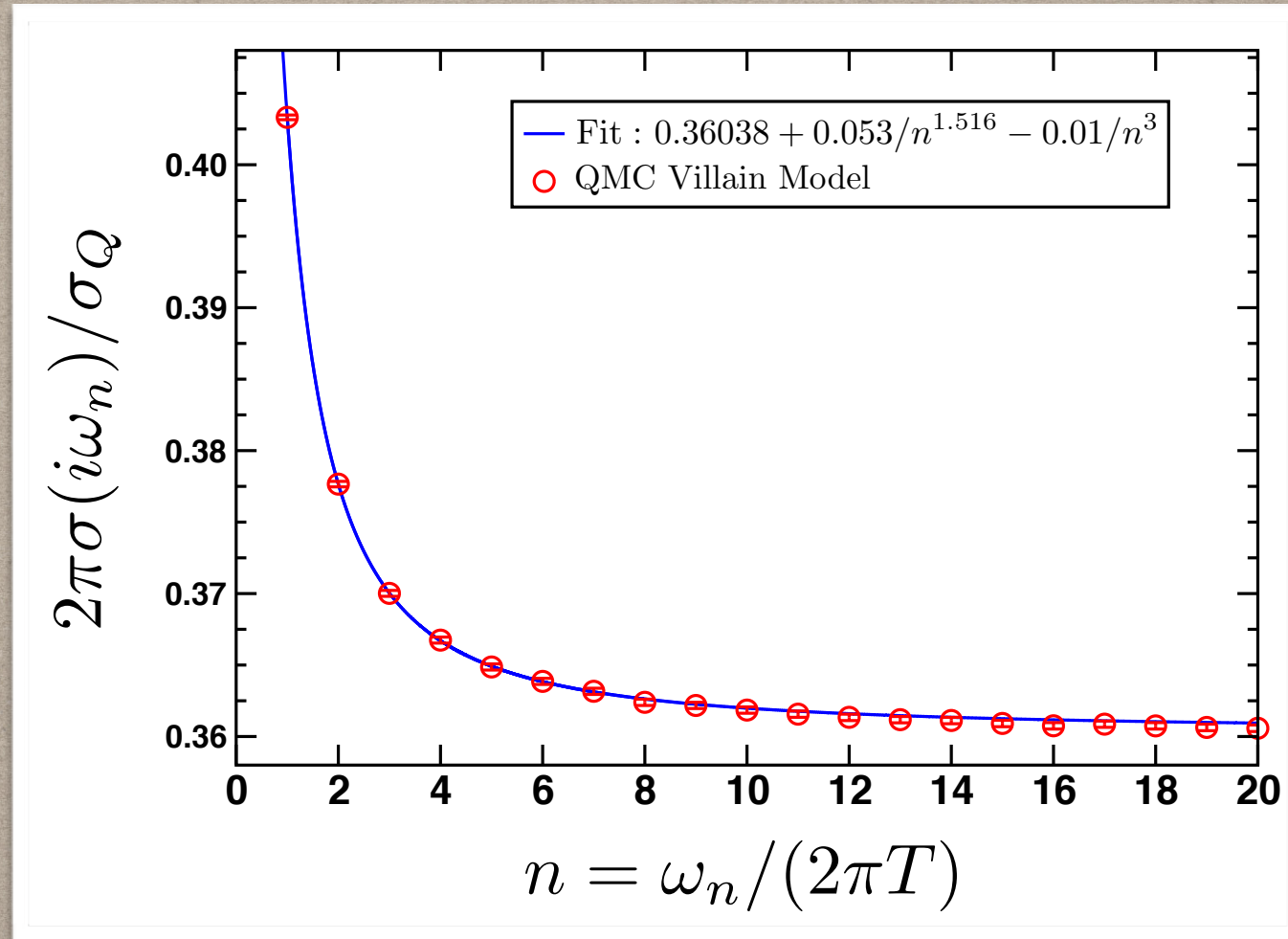
$$\lim_{|\omega_n| \gg p} J_x(\omega) J_x(-\omega + p) = -|\omega_n| \sigma_\infty \delta^{(3)}(p) - \frac{\mathcal{C}}{|\omega_n|^{\Delta-1}} \mathcal{O}(p) \\ + \frac{\mathcal{C}_T}{\omega_n^2} \left[T_{xx}(p) - T_{yy}(p) - 12\gamma(T_{xx}(p) + T_{yy}(p)) \right] + \dots$$



$$\frac{\sigma(i\omega_n)}{\sigma_Q} = \sigma_\infty + b_1 \left(\frac{T}{\omega_n} \right)^\Delta + b_2 \left(\frac{T}{\omega_n} \right)^3 + \dots, \quad \omega_n \gg T$$

Thermal Operator, \mathcal{O}_g scaling dimension $\Delta = 3 - 1/\nu$

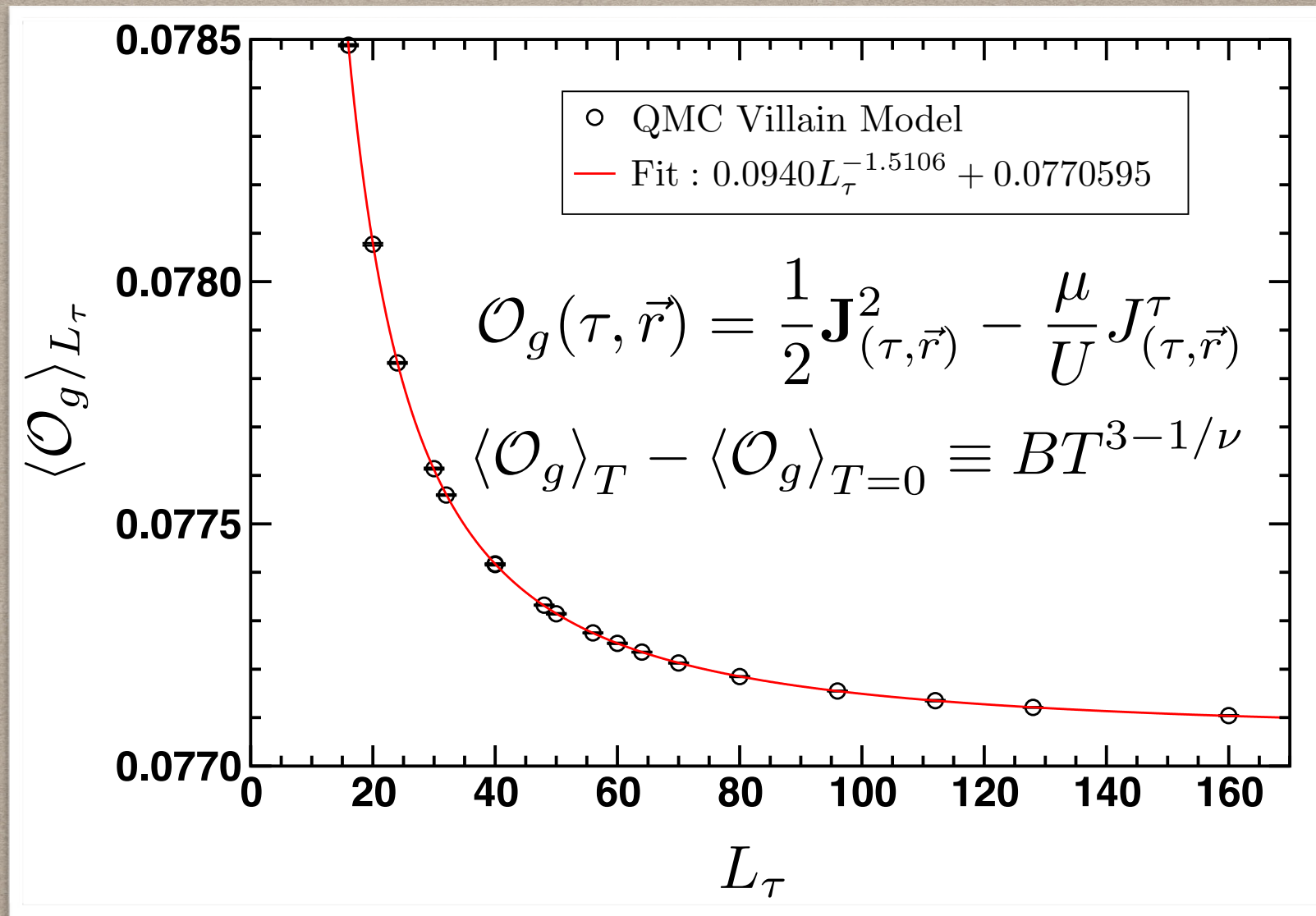
E.Katz, S. Sachdev, ESS, W. Witczak-Krempa Arxiv (2014)



$$\sigma_\infty/\sigma_Q = 0.36038\dots$$

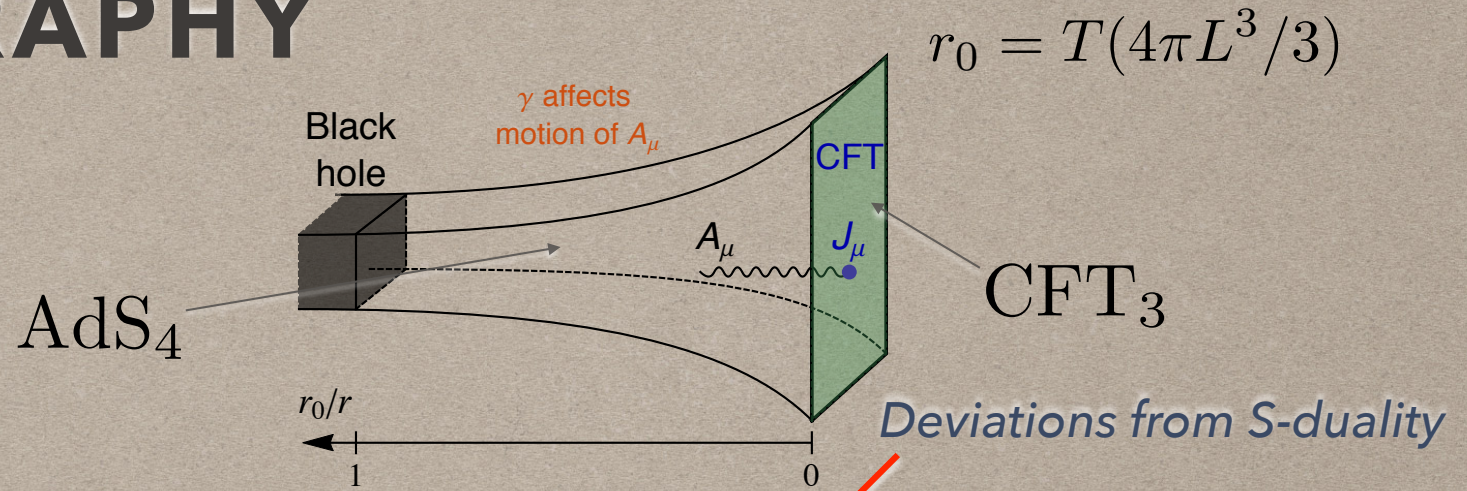
$$\nu = 0.674$$

EXISTENCE OF THERMAL OPERATOR



$$\nu = 0.6714(10)$$

HOLOGRAPHY



$$S_{\text{bulk}} = \int d^4x \sqrt{-g} \left[\frac{1}{8\pi G} \left(R + \frac{6}{L^2} \right) - \frac{1}{4g_4^2} F_{ab} F^{ab} + \gamma \frac{L^2}{g_4^2} C_{abcd} F^{ab} F^{cd} \right] \quad |\gamma| \leq \frac{1}{12}$$

R , Riemann curv., L Radius of Curvature of AdS₄, C Weyl curvature tensor

$$\sigma(w) = -\frac{i}{3w} \frac{\partial_u A_y}{A_y} \Big|_{u=0}$$

$$\sigma \left(\frac{\omega}{T} \rightarrow \infty \right) = 1/g_4^2$$

2 Parameter fit in $\sigma(\infty)$ and γ

$$A_y'' + \left(\frac{f'}{f} + \frac{g'}{g} \right) A_y' + \frac{9w^2}{f^2} A_y = 0$$

$$f(u) = 1 - u^3, \quad g(u) = 1 + 4\gamma u^3, \quad u = r_0/r$$

Modified Maxwell equation

R. Myers, S. Sachdev, A. Singh, PRD (2011)

W. Witczak-Krempa, S. Sachdev, PRB (2013)

REFINED ANALYSIS: INCLUDE SCALAR FIELD

$$S = \int d^4x \sqrt{-g_{\text{Sch}}} \left[\frac{1}{\kappa^2} \left(R + \frac{6}{L^2} \right) - \frac{1}{4g_4^2} [1 + \alpha\phi(u)] F_{ab} F^{ab} \right]$$

$$8\pi G$$

$$\sigma \left(\frac{\omega}{T} \rightarrow \infty \right) = 1/g_4^2$$

$$b_1 = \sigma_\infty \alpha a \frac{\Gamma(\Delta + 1)}{2^\Delta} \left(\frac{4\pi}{3} \right)^\Delta$$

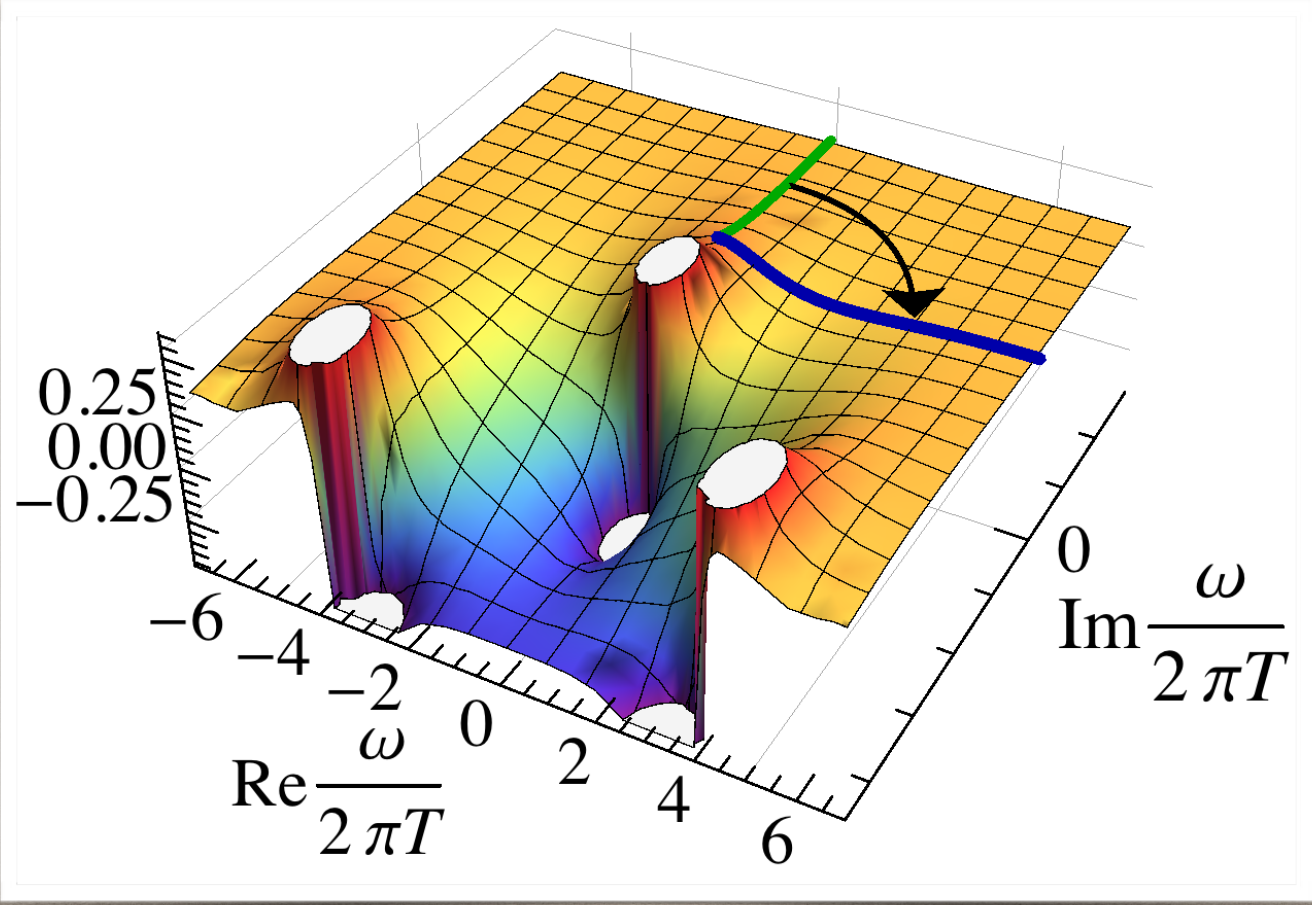
$$\Delta_g = 3 - 1/\nu$$

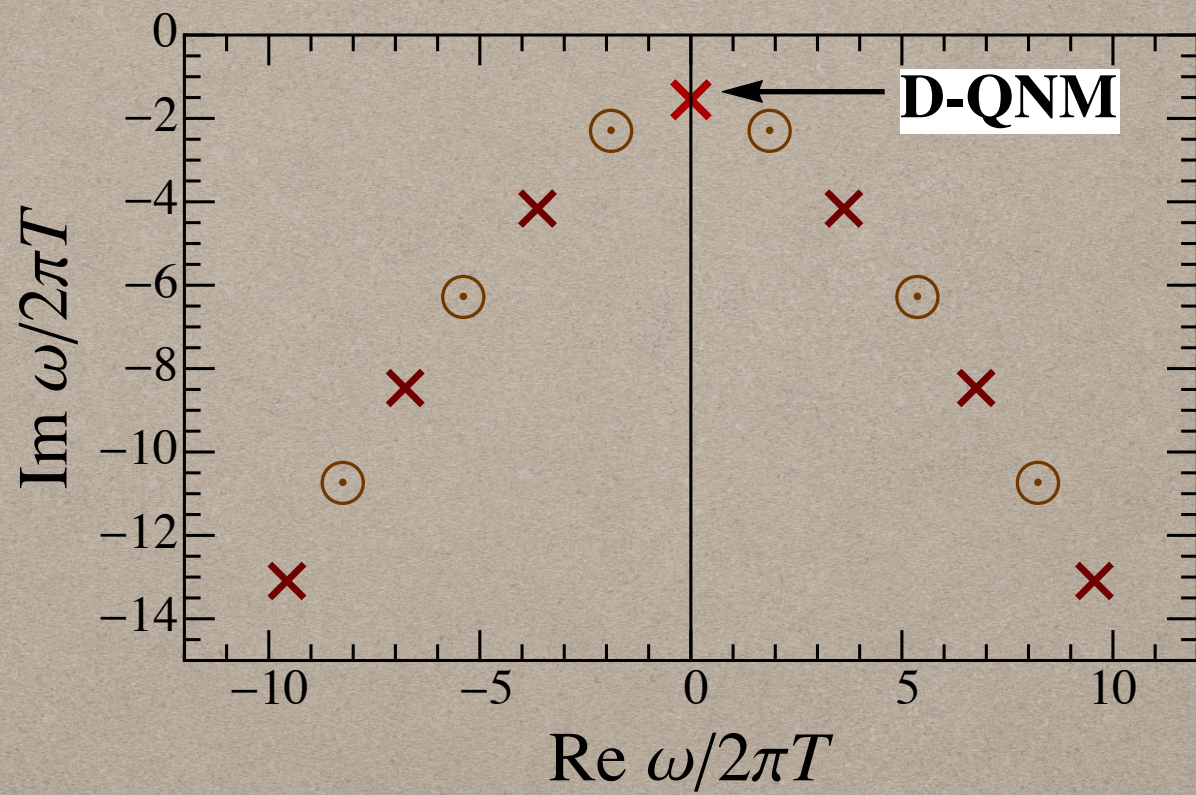
$$\frac{\sigma(i\omega_n)}{\sigma_Q} = \sigma_\infty + b_1 \left(\frac{T}{\omega_n} \right)^\Delta + \dots$$

$$\left((1 + \alpha\phi) f A'_y \right)' - w_n^2 \frac{(1 + \alpha\phi)}{f} A_y = 0 \quad ; \quad w_n \equiv \frac{3\omega_n}{4\pi T}$$

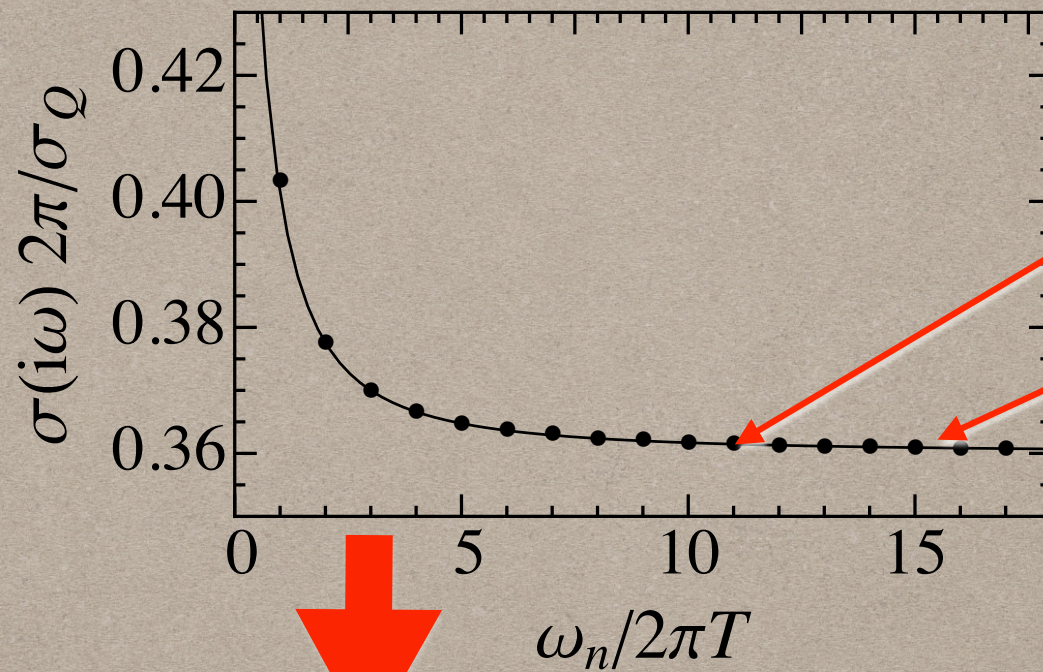
2 Parameter fit in $\sigma(\infty)$ and b_1

HOLOGRAPHIC CONDUCTIVITY IN THE COMPLEX PLANE

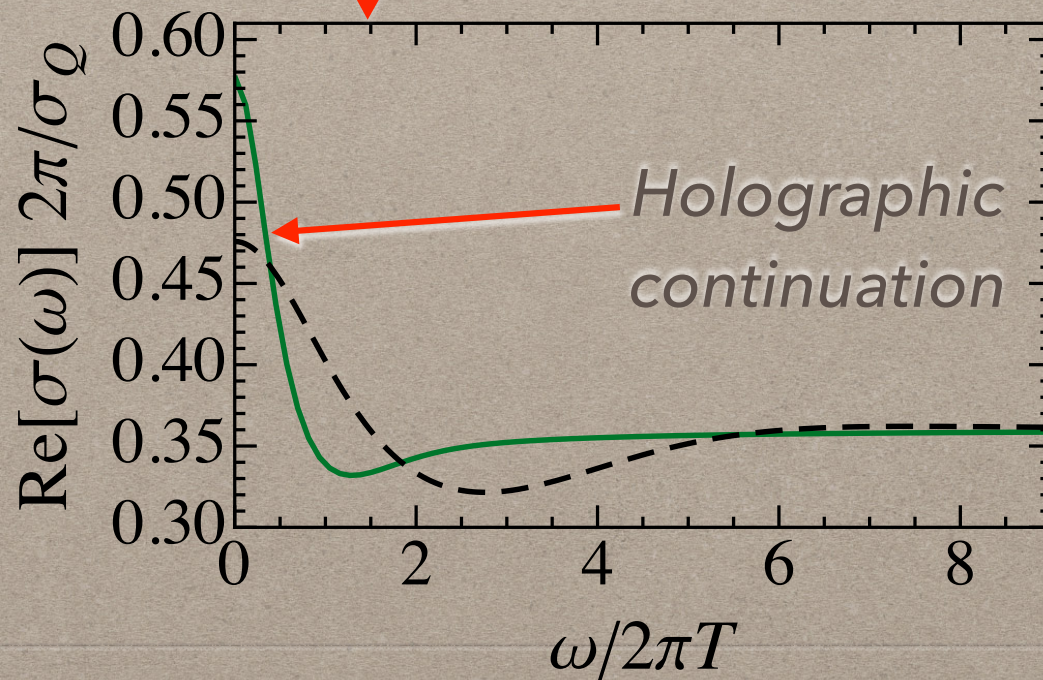




Imag Axis



Real Axis



CONCLUSIONS

- * Particle like excitations
- * Sum rules

$$\int_0^\infty d\omega [\text{Re } \sigma(\omega/T) - \sigma(\infty)] = 0$$

- * New QMC results for Universal $\sigma(i\omega/T)$ using 2 different models
- * AdS/CFT continuation

